

Identification of Broadband Source-Array Responses from Sensor Second Order Statistics

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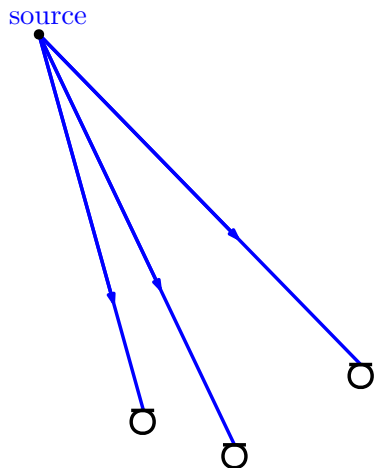
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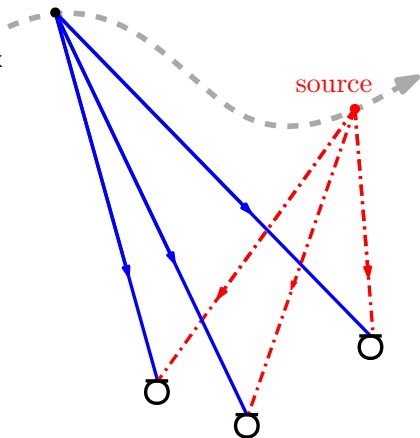
Presentation Overview

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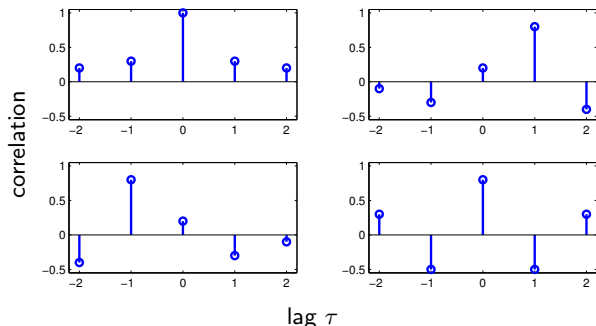
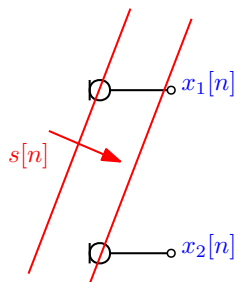
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Space-Time Covariance

- ▶ We have M sensor signals organised in $\mathbf{x}[n] \in \mathbb{C}^M$;
- ▶ to take the broadband nature of signals into account, we must consider lags τ ;
- ▶ space-time covariance matrix $\mathbf{R}[\tau] = \mathcal{E}\{\mathbf{x}[n]\mathbf{x}^H[n - \tau]\}$;



Cross-Spectral Density Matrix

- ▶ CSD matrix forms a z -transform pair with the space-time covariance matrix,

$$\mathbf{R}(z) = \sum_{\tau} \mathbf{R}[\tau] z^{-\tau} \quad \text{or} \quad \mathbf{R}(z) \bullet \text{---} \circ \mathbf{R}[\tau] ;$$

- ▶ symmetry of $\mathbf{R}[\tau] \longrightarrow \mathbf{R}(z)$ is parahermitian:

$$\mathbf{R}(z) = \mathbf{R}^P(z) = \mathbf{R}^H(1/z^*) ;$$

(Hermitian transposition and time reversal)

- ▶ link to a narrowband covariance at normalised angular freq. Ω_k ,

$$\mathbf{R}(e^{j\Omega_k}) = \mathbf{R}(z) \Big|_{z=e^{j\Omega_k}}$$

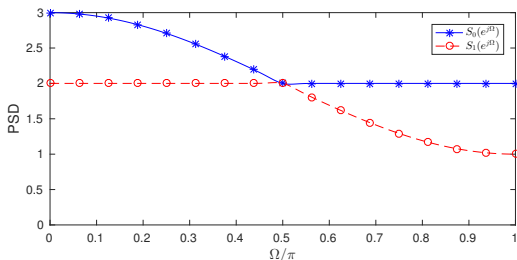
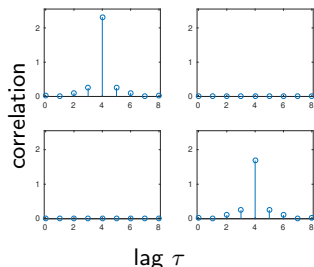
- ▶ many optimal (narrowband!) methods are based on decompositions such as the EVD: $\mathbf{R}(e^{j\Omega_k}) = \mathbf{Q}\mathbf{\Lambda}\mathbf{Q}^H$.

McWhirter Decomposition

- ▶ John McWhirter et al. (2007): polynomial eigenvalue decomposition of a parahermitian matrix:

$$\mathbf{R}(z) \approx \mathbf{Q}(z)\mathbf{\Lambda}(z)\mathbf{Q}^P(z)$$

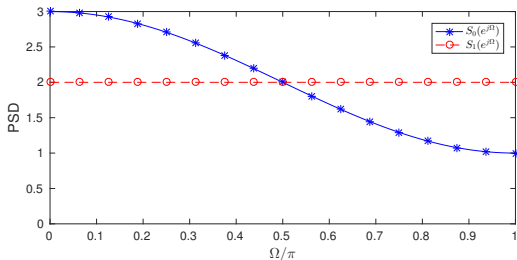
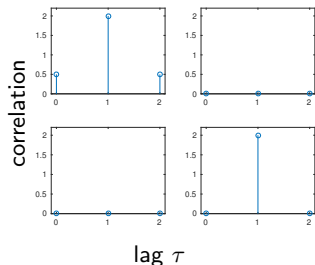
- ▶ paraunitary (i.e. lossless) matrix $\mathbf{Q}(z)$, s.t. $\mathbf{Q}(z)\mathbf{Q}^P(z) = \mathbf{I}$;
- ▶ diagonal and spectrally majorised $\mathbf{\Lambda}(z)$:



Parahermitian Matrix EVD (PEVD)

- ▶ Franz Rellich (1937): for $\mathbf{R}(e^{j\Omega})$ analytic, there exist analytic eigenvectors $\mathbf{\Gamma}(e^{j\Omega})$ and analytic eigenvalues $\mathbf{U}(e^{j\Omega})$;
- ▶ can be generalised to

$$\mathbf{R}(z) = \mathbf{U}(z)\mathbf{\Gamma}(z)\mathbf{U}^P(z) ;$$

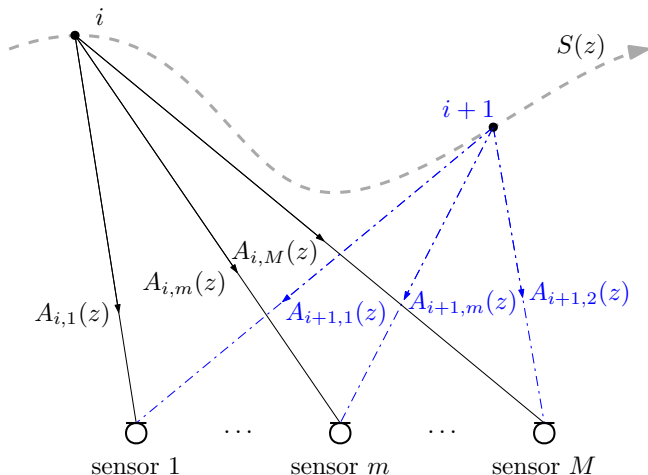


- ▶ eigenvalues are unique, eigenvectors can be modified by arbitrary allpass filters $H(z)$ (s.t. $H(z)H^P(z) = 1$),

$$\mathbf{R}(z)\mathbf{u}(z)H(z) = \gamma(z)\mathbf{u}(z)H(z) .$$

Source-Sensor Transfer Functions

- ▶ We take M -array measurements of a single source:



- ▶ 2nd order stats: $\mathbf{R}_i(z) = S(z)\mathbf{a}_i(z)\mathbf{a}_i^P(z) = \gamma_i(z)\mathbf{u}_i(z)\mathbf{u}_i^P$.

Transfer Functions and PEVD

- ▶ 2nd order stats: $\mathbf{R}_i(z) = S(z)\mathbf{a}_i(z)\mathbf{a}_i^P(z) = \gamma_{i,m}(z)\mathbf{u}_i(z)\mathbf{u}_i^P$;
- ▶ difference: $\mathbf{u}_i(z)$ is normal, $\mathbf{u}_i^P(z)\mathbf{u}_i(z) = 1$, while $\mathbf{a}_i(z)$ is not:

$$\mathbf{a}_i^P(z)\mathbf{a}_i(z) = A_{i,(-)}(z)A_{i,(+)}(z) = A_{i,(+)}^P(z)A_{i,(+)}(z)$$

with minimum-phase $A_{(+)}(z)$;

- ▶ therefore:

$$H_i(z)\mathbf{u}_i(z) = \frac{\mathbf{a}_i(z)}{A_{i,(+)}(z)}$$

$$\gamma_i(z) = A_{i,(+)}(z)S(z)A_{i,(+)}^P(z) ,$$

- ▶ from a single measurement $\mathbf{R}_i(z)$, we cannot say anything about $\mathbf{a}_i(z)$ or $S(z)$.

Multiple Measurements

- ▶ If we have several measurements $i = 1 \dots I$:

$$H_i(z)\mathbf{u}_i(z) = \frac{\mathbf{a}_i(z)}{A_{i,(+)}(z)}$$
$$\gamma_i(z) = A_{i,(+)}(z)S(z)A_{i,(+)}^P(z),$$

- ▶ we can extract $S(z)$ as the greatest common divisor

$$\hat{S}(z) = \text{GCD}\{\gamma_1(z) \dots \gamma_I(z)\};$$

- ▶ we can then extract the terms $A_{i,(+)}(z)$, and hence determine the vectors $\mathbf{a}_i(z)$ save of an arbitrary phase response due to the allpass $H_i(z)$:

$$\mathbf{a}_i(z) = A_{i,(+)}(z)H_i(z)\mathbf{u}_i(z).$$

Alternative DFT Domain Attempt

- ▶ As an alternative, we take measurements in independent frequency bins $\Omega_k = \frac{2\pi k}{K}$:

$$\begin{aligned} \mathbf{R}_{i,k} &= \mathbf{R}_i(e^{j\Omega_k}) = \mathbf{a}_i(e^{j\Omega_k})S(e^{j\Omega_k})\mathbf{a}_i^H(e^{j\Omega_k}) \\ &= \mathbf{q}_{i,k}\lambda_{i,k}\mathbf{q}_{i,k}^H. \end{aligned}$$

- ▶ principal eigenvectors and eigenvalues for the measurement campaigns are

$$\begin{aligned} \mathbf{q}_{i,k} &= \frac{\mathbf{a}_i(e^{j\Omega_k})}{|\mathbf{a}_i(e^{j\Omega_k})|}, \\ \lambda_{i,k} &= S(e^{j\Omega_k})|\mathbf{a}_i(e^{j\Omega_k})|^2. \end{aligned}$$

- ▶ because of the normalisation, nothing can be extracted about the source or the transfer functions.

Numerical Example

- ▶ Source with power spectral density

$$S(z) = \frac{1}{2}z + \frac{5}{4} + \frac{1}{2}z^{-1}$$

- ▶ vector of transfer functions during campaign $i = 1$:

$$\mathbf{a}_1(z) = \begin{bmatrix} 1 & + & \frac{1}{2}z^{-1} \\ \frac{3}{4} & - & \frac{1}{2}z^{-1} \end{bmatrix}$$

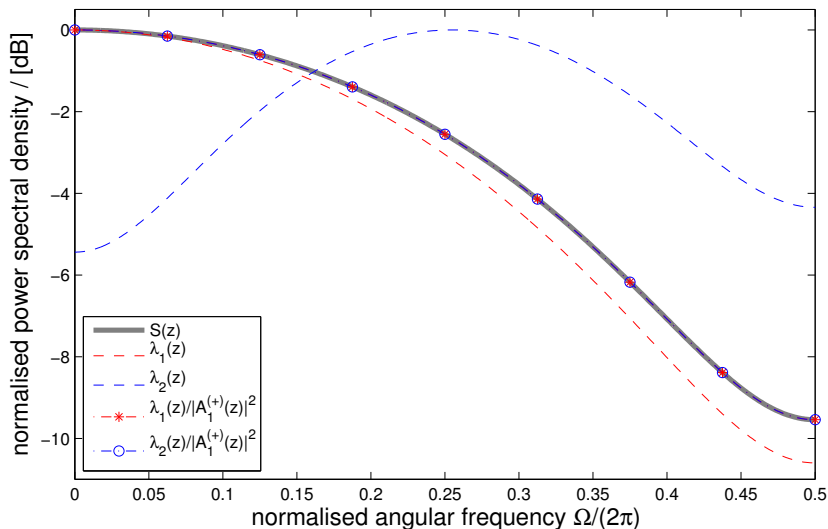
- ▶ vector of transfer functions during campaign $i = 2$:

$$\mathbf{a}_2(z) = \begin{bmatrix} \frac{4}{5} & - & \frac{1}{2}z^{-1} \\ -\frac{1}{2} & + & z^{-1} \end{bmatrix} ;$$

- ▶ based on these: PEVD computations for $\mathbf{R}_1(z)$ and $\mathbf{R}_2(z)$, and GCD calculation based on eigenvalues.

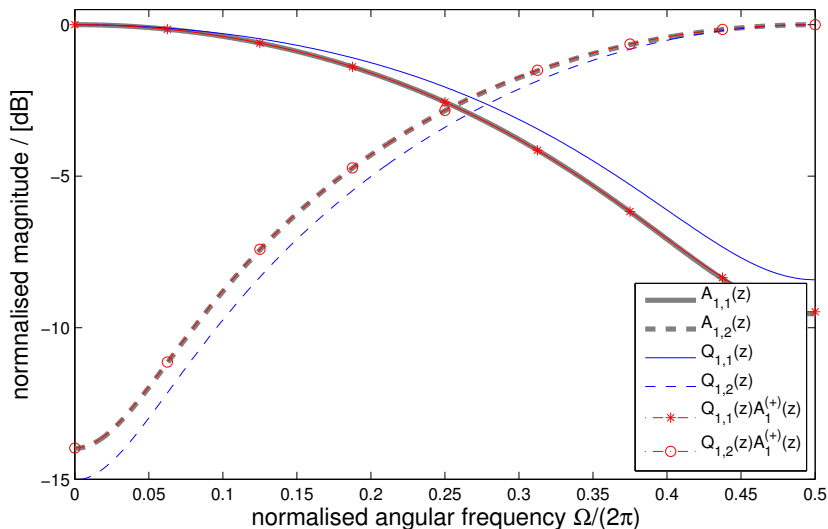
Numerical Results — Source PSD

- ▶ Eigenvalues / source PSD for both measurements $i = \{1, 2\}$:



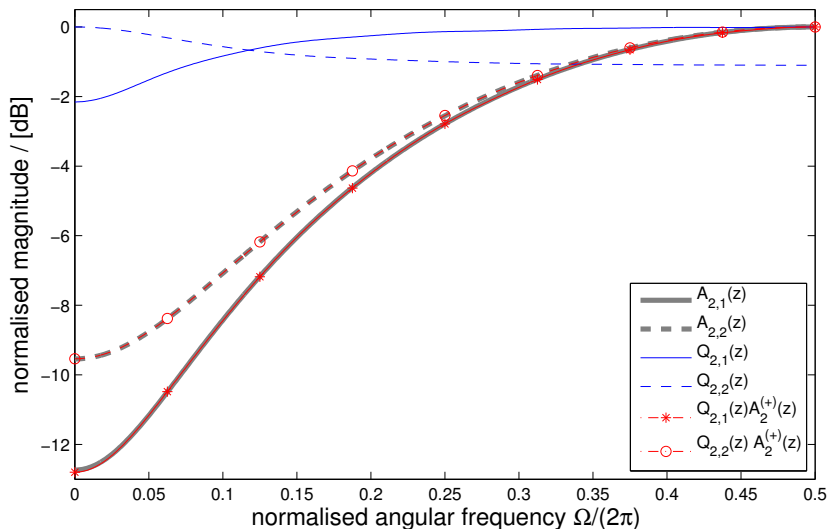
Numerical Result — Magnitude Responses I

- Eigenvectors / magnitude response for measurement $i = \{1\}$:



Numerical Result — Magnitude Responses II

- Eigenvectors / magnitude response for measurement $i = \{2\}$:



Summary and Critique

- ▶ We can extract the source PSD and the magnitude responses once we have at least two measurements;
- ▶ an independent frequency bin approach does not yield anything;
- ▶ the polynomial approach rests on an accurate parahermitian EVD, and an accurate root finding / GCD algorithm;
- ▶ root finding is numerically challenging: research since Euclid (300BC), with robust root-finding methods still on-going (Gröbner bases, algebraic geometry);
- ▶ nevertheless the approach gives a glimpse of the type of advantages that a coherent broadband approach can offer.